Radiative de-excitation channel of slow highly charged ions transmitted through freestanding single layer graphene

R.A. Wilhelm^{*1}, J. Schwestka[†], E. Gruber[†], R. Heller^{*}, R. Kozubek[‡], M. Schleberger[‡], S. Facsko^{*}, F. Aumayr[†]

* Helmholtz-Zentrum Dresden-Rossendorf, Institute of Ion Beam Physics and Materials Research, 01328 Dresden,

Germany

[†] TU Wien, Institute of Applied Physics, 1040 Vienna, Austria

[‡] University Duisburg-Essen, Faculty of Physics and CENIDE, 47057 Duisburg, Germany

Synopsis X-ray emission spectra are recorded from slow (<1 a.u.) highly charged ions with and without K-shell vacancies in coincidence with charge exchange and energy loss measurements of these ions after they have passed a single layer of graphene.

Slow highly charged ions interacting with a solid surface undergo an ultrafast charge exchange combined with a rapid electronic de-excitation within less than 10 fs [1, 2]. These processes involve capture of some 10 electrons, emission of at least some 10 additional electrons from the surface and radiative as well as non-radiative de-excitation of the ion. To investigate the branching ratio of radiative vs. nonradiative de-excitation we measured x-ray emission of highly charged Ar ions with 2, 1 or no K-shell hole(s) when they are transmitted through a freestanding single layer of graphene. To discriminate x-ray emission from ions impacting on the sample holder or the graphene support grid structure we use a coincidence technique where ions are measured after transmission in a channeltron. Additionally the channeltron is mounted at the exit of an electrostatic analyzer, thus we have also access to the transmitted ion's energy and charge state.

Fig. 1 and 2 show x-ray emission spectra from Ar^{17+} and Ar^{18+} , respectively, due to their interaction with graphene.



Figure 1. Ar¹⁷⁺ x-ray emission due to the ion interaction with graphene. The data is compared to calculations by Bhalla [3]

Peaks between the K_{α} - and the Ly_{β} -line are observed. These peak positions and their ratios can very

¹E-mail: r.wilhelm@hzdr.de

well be reproduced by comparison to atomic structure code calculations by Bhalla and Mirakhmedov [3, 4] when their values and ratios are convoluted with a Gauss shaped detector function with a FWHM of 140 eV (energy resolution of the Bruker x-flash silicon drift detector).



Figure 2. Ar^{18+} spectrum as in Fig. 1. The data is compared to calculations by Mirakhmedov [4]

From this comparison we see that the radiative Kshell emission occurs while the outer shells are not yet fully filled. Additionally, since we use a windowless silicon drift detector we are also sensitive to emitted electrons. From the spectrum we can directly identify Ar KLL Auger electrons (after passing through a thin dead layer in the detector surface) and estimate the branching ratio between radiative and non-radiative de-excitation for our experimental situation.

References

- [1] E. Gruber et al. 2016 Nature Commun. 7 13948
- [2] R.A. Wilhelm et al. 2015 Prog. Surf. Sci. 90 377–395
- [3] C.P. Bhalla 1973 Phys. Rev. A 8 2877
- [4] M.N. Mirakhmedov 1995 *Nucl. Instrum. Meth. B* 98 429–435