Shubnikov-de Haas oscillations in AIN/GaN/AIN quantum-wells on single-crystal AIN substrates

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AlN/GaN/AlN quantum well high electron mobility transistor (QW HEMT) heterostructures feature large energy band offsets between GaN well of ≈ 20 nm surrounded by AlN back barrier of ≈ 500 nm and AlN top barrier of ≈ 5 nm. The AlN layers have no strain as they are grown on single-crystal bulk substrate, whereas the GaN channel is under a large compressive strain of -2.3%, which modifies its energy-momentum dispersion. The large discontinuity in spontaneous and piezoelectric polarization between GaN and AlN induces a 2D electron gas (2DEG) of high density (~2x10¹³ cm²) with highly populated electronic sub-bands. This undoped binary heterostructure typically exhibits electron mobilities of approximately 1200 cm²/Vs at 10 K, limited by Stark-effect scattering from a) the strong internal electric field in the well, and b) Coulomb drag, and parallel conduction between the 2DEG and polarization-induced 2D hole gas (2DHG) present on opposite sides of the well forming a bilayer in undoped control samples [Figure 1 (a)].

To remove the 2DHG and reduce the electric field in the well, we incorporated n-type compensation δ -doping in the AlN/GaN/AlN heterostructure [Figure. 1 (b)], which resulted in enhanced Hall-mobilities of 854.1 cm²/Vs at RT and 2240.2 cm²/Vs at 10 K, and as a result record low sheet resistances (86.9 Ω/\Box at 10 K) [Table. 1]. Shubnikov-de Haas (SdH) oscillations in the longitudinal magnetoresistance were observed in both heterostructures with an onset of ~ 8 Tesla [Figures. 2(a) and 2(b)]. The Fast Fourier Transform (FFT) of ΔR_{xx} reveals single oscillation frequency, indicating one sub-band occupation [inset of Figs. 2(c) and 2(d)]. Analysis of the FTT amplitudes indicates an electron effective mass m^{*} \approx 0.26 m₀ for the un-doped QW and \approx 0.28 m₀ for the δ -doped QW [Figures. 2(c) and 2(d)], slightly higher than \approx 0.2-0.23 m₀ in conventional AlGaN/GaN heterostructures of lower 2DEG densities, possibly due to the compressive strain of the GaN channel and the strong nonparabolicity of the sub-band at high energies. Finally, δ -doped QW HEMTs exhibit twice the quantum scattering lifetime τ_q than the un-doped counterpart. A Dingle ratio $\tau_{classic}/\tau_q \cong 2.58$ at 2 K suggests the prevalence of short-range scattering potentials, likely arising from interface roughness (IR) scattering [Figures. 2 (e)-(f)].



Un-doped Un-doped Un-doped (a) (c) (e) 10 0.6 40 0.4 (a.u.) ARXX (Ohms) 30 ARXX (Ohms) 0.2 amplitude FFT 0.0 10 200 400 600 80010 20 Frequency (Tesla) -0.2 m*: 0.255 m₀ FFT 10 $\tau_a = 71$ fs @ 2K -0.4 $n_s = 1.8 \times 10^{13} \text{ cm}^{-2}$ 0 L 0 10 -0.6 -0.07 0.075 0.080 0.085 0.090 0.09 0.10 10 20 0.08 0.11 30 1/Magnetic Field (1/Tesla) 1/Magnetic Field (1/Tesla) Temperature (K) (b) (d) (f) δ-doped δ-doped δ-doped 0.6 10 40 0.4 amplitude (a.u.) ARXX (Ohms) ARxx (Ohms) 0.2 30 H 0.0 10 20 200 400 600 800 10 equency (Tesla) -0.2 FFT m*: 0.280 m₀ 10 $\tau_a = 133$ fs @ 2K -0.4 $= 3.1 \times 10^{13} \text{ cm}^{-3}$ -0.6 -0.07 0 L 10 0.075 0.080 0.085 0.090 20 0.08 0.09 0.10 0.11 10 30 Temperature (K) 1/Magnetic Field (1/Tesla) 1/Magnetic Field (1/Tesla)

Figure 1. Schematics of the (a) un-doped and (b) δ -doped AlN/GaN/AlN heterostructures. Table 1. 2DEG densities, mobilities and sheet resistances measured via Hall-effect at 10 K.

